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extraordinary thermal and photochemical stability of widebandgap oxide semiconductors lends them to extensive application for light conversion, but a potential drawback with these materials is their large bandgap which requires high energy light to create electron-hole pairs. The problem has largely been solved by sensitizing them to visible light with dyes. Unfortunately, the light-to-electricity conversion harvesting efficiency with only one layer of absorbent dye or low charge separation and transfer efficiency with thick dye layers on a flat surface electrode.5-7 Recently several improvements have substantially raised the light-to-electricity conversion efficiency of dye-sensitized solar cells:8-26 (1) high-፤。 flat electrodes have been used in order to increase incident light harvesting efficiency through more dye being attached ፤ de lectrode, which de lectrode efficiency with one layer dye adsorption on a flat electrode and low charge injection with thick dye. (2) Improvement of adsorption performance between dye and electrode using a dye containing strong adsorbent groups which make the dye ढ<table-cell> disadvantage of short excited state lifetime of the dye and of more efficient sensitizers. Extensive investigations have shown that porphyrins and complexes of second- and thirdrow transition metals with polypyridyl ligands possess in particular, have proven to be the most efficient sensitizers. Gratzel and coworkers in particular have pioneered this field. $^{8\text{--}18}$ They obtained $>\!10\%$ light-to-electricity conversion These results make the practical application of photoelectrochemical cells feasible.^{15,22} In order to further improve sensitization of dyes and conversion efficiency of solar cells to

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Scheme 1 Structures of ligands used in here and complexes 1, 2, 3

Experimental

Chemicals

Syntheses

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Electrode preparation

Transparent TiO₂ films were prepared from commercial TiO₂ powder (Degussa, P25) in a manner analogous to that previously described:^{15,23} 6 g TiO₂ powder were ground in a marble mortar with 2 ml water containing 0.2 ml acetylacetone to prevent reaggregation of the particles. After the powder had been dispersed by a high shear force in the viscous paste, it was diluted by slow addition of water (8 ml) under continued facilitate the spreading of the colloid on the substrate. A 3×6 cm plate 30Ω \Box^{-1} conducting glass was used as the substrate for the deposition of the TiO₂. The plate was scribed down the middle and at 1 cm intervals along its length (6 cm). A cellophane tape of 0.5 cm width was glued on each edge of of the TiO_2 layer as well as to mask an area of the electrode TiO₂ solution were distributed on the plate surface and spread uniformly by rolling a 1 cm diameter test tube across the substrate surface, and allowed to air dry. The substrate coated TiO_2 then was heated in a furnace in which the temperature was increased gradually to 500 °C and then kept at 500 °C for 30 min. After this, the substrate was broken into 1.5×1 cm pieces and 1 cm^2 active surface area TiO₂ films were obtained.

Surface attachment of complexes

Electrochemistry

 working electrode was manually cleaned and polished prior to each individual scan.

Optical measurements

Absorption spectra were recorded using two matched 1 cm glass cells on a Shimadzu UV-1600A recording spectrophotometer. Samples were prepared gravimetrically as $ca. 3 \times 10^{-5}$ M ethanol solutions. IR spectra were recorded on KBr pellets with a Perkin Elmer 983G spectrometer. Emission spectra were recorded on a Hitachi 850 fluorescence spectrometer. Emission spectra were collected at maximal absorption in the visible region as excitation wavelength and reported as uncorrected values. Quantum yields were obtained using $[Ru(bpy)_3][PF_6]_2$ in MeCN as a quantum counter calibrant. Emission lifetimes were measured by nanosecond laser flash photolysis using a DCR Nd:YAG laser (532 nm) as the excitation source and SMA (spectrometric multichannel analyzer, Princeton Instruments Inc. OSMA Detector Controller) as detector. The data were then transferred to a 486PC to calculate the lifetimes of the complexes. Samples for emission lifetime and quantum yield measurements were saturated with argon gas for 20 min.

Photoelectrochemistry

Results and Discussion

Photophysical properties

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Fig. 1 Absorption of complexes 1, 2, 3 in ethanol at room temperature

followed by 2 then 1. This is attributed to π back-bonding, *i.e.*, electron donation from Ru^{II} t_{2g} orbitals to empty π^* levels of the ligands, consistent with the π^* -accepting ability of biquin > phen > dmpy, which is reflected by their free ligand reduction potentials.

Fig. 2 shows the absorption spectra of complexes 1, 2 and 3 adsorbed onto TiO₂ films as well as in ethanol solution. Comparing their absorption in films with in solution, we found that the profiles of the absorption spectra of 1 and 2 in films changed significantly showing new absorption bands in the red region. However for 3, no new absorptions were observed. This can be explained as follows: the interaction of the phosphonate group with surface Ti ions is likely via formation of P-O-Ti bonds. Thus, the role of the PO_3H_2 is to serve as an interlocking group coupling electronically the π^* orbitals of the terpy to the Ti 3d orbital manifold of the semiconductor similarly to COOH.¹⁴ This coupling increases the delocalization of the π^* orbitals of the terpyridine resulting in a decrease of the energy of the π^* level and the red shift of absorption of 1 and 2 in the TiO₂ film. For 3, the energy of the π^* level of terpyridine is also lowered as for 1 and 2, therefore, it might have been expected to show a redshifted absorption spectrum. However, it is known, for polypyridyl ruthenium(II) complexes, that the excited electron is localized on a single ligand, in complexes containing a mixed ligand set.33 The excited electron is always localized on the ligand which is easier to reduce.³³ Biquin is more easily reduced than terpyridine, and the energy of the π^* orbital of biquin is much lower than that of



Fig. 2 Absorption spectra of complexes 1, 2, 3 on the TiO_2 film electrode, using a bare TiO_2 film electrode as a reference. — in film electrode, --- in ethanol.

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complex	s_	pH 3	pH 9	ণ্ণ	S S S S S S S S S S S S S S S S S S S S	革	$E_{1/2}{}^{\rm ox}/{\rm V}^d$
1	%%	490	506	هـــــــــــــــــــــــــــــــــــــــ	উ	近	0.86
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3	उ	561	581	줞цининининининининининининининининининин	ජ	ذ	1.00

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Table 2 Calculated LUMO energies of different ligands^a

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^aThe two pyridine rings in dmpy and the three pyridine rings in ttp were constrained to be planar.

terpyridine (Table 2). The LUMO energies of ligands calculated by AM1 semi-quantitative quantum calculations reflect the ability of the ligands to accept an electron.

Owing to the large LUMO energy difference between biquin and ttp, even if the π^* orbital energy (or LUMO) of the latter is somewhat decreased when attached onto TiO₂ electrode, this is not enough to alter the sequence of the π^* orbital energy levels. Thus the excited electron is localized on biquin in complex 3 whether it is in solution or in TiO_2 film. The interaction of 3 with the TiO₂ surface does not involve the ligand relevant to MLCT, so the electronic spectrum of 3 adsorbed onto TiO₂ electrode does not show an observable change compared with its spectrum in solution. For 2, while the LUMO energy of phen is slightly lower than that of ttp, the decrease of π^* orbital energy (or LUMO) of ttp caused by the interaction with TiO₂ results in switching of their level order. When 2 is adsorbed onto TiO₂, the excited electron is now localized on ttp. The LUMO energy of dmpy in complex 1 is much higher than that of ttp so the excited electron is mainly localized on ttp both in solution and in the TiO₂ film. The interaction of either 1 or 2 with TiO₂ involves the MLCT ligand, so complexes 1 and 2 attached onto TiO_2 film electrodes exhibit electronic spectra different from their spectra in solution.

Photoelectrochemical properties

Fig. 3 shows a plot of incident-photo-to-current-efficiency, IPCE(λ), for three metal complexes anchored to TiO₂ electrodes. The measurements were made with an electrometer in a two-electrode arrangement and a 0.3 M LiI–0.03 M I₂ propylene carbonate electrolyte solution. The IPCE(%) is defined by eqn. (1):

IPCE(%) =

$$\frac{1.25 \times 10^3 \text{ (eV nm)} \times \text{photocurrent density } (\mu \text{A cm}^{-2})}{\text{wavelength (nm)} \times \text{photoflux } (\mu \text{W cm}^{-2})}$$
(1)

Fig. 3 demonstrates that the IPCE is dependent on the surface attached compound. Maximum observed IPCE follows the trend: $1 > 2 \gg 3$ with complex 3 attached to TiO₂ film electrode showing virtually no photosensitization. The maximal IPCE of the TiO₂ film cell sensitized by 1 was >75% cf. 45% for 2. The performances of the photoelectrochemical cells under white light illumination are is given in Table 3. Short circuit photocurrent densities, I_{sc} , and open circuit photovoltages V_{oc} were measured as described in the Experimental section. Current–voltage curves are shown in Fig. 4.



Fig. 3 Photoaction spectra of the sensitizers bound to ${\rm TiO}_2$ in a regenerative solar cell

The fill factors were calculated from current–voltage curves. The best performance was shown by complex **1**. A key parameter of cell performance is the incident photon to current efficiency (IPCE), which directly reflects how efficiently incident photons lead to electrons. The IPCE is composed of three terms, expressed by eqn. (2):

$$IPCE = (LHE)\phi\eta \tag{2}$$

where LHE is the light-harvesting efficiency, ϕ is the quantum yield for charge injection and η is the efficiency of collecting electrons on the external circuit, which reflects the probability that the injected electrons escape from recombination. In order to delineate the respective influences of the three factors on IPCE, it is necessary to compare the three parameters in eqn. (2).

Light harvesting efficiency

For a ideal sensitizer, all incident light power would be absorbed (LHE=1), but this is actually never observed. The absorption factor (α) is the product of the radiant power absorbed by the system, *I*, divided by the incident radiant power, I_0 , $\alpha = I/I_0$.³⁴ In photoelectrochemical literature, the absorption coefficient is expressed as the LHE. If losses are caused only by transmission of light, LHE = $\alpha = I/I_0 = 1 - T =$ $1 - 10^{-4}$, where *T* is the transmittance and *A* is the absorbance.³⁴ Assuming Beer's law is applicable, LHE (λ) can be related to the molar absorption coefficient by eqn. (3):

$$LHE(\lambda) = 1 - 10^{-\epsilon\Gamma}$$
(3)

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Fig. 4 Photocurrent-curves of the photoanodes in a regenerative dyesensitized solar cell under white light excitation. 0.2 cm^2 active area was radiated.



Fig. 5 An accepted model for dye sensitization in regenerative photoelectrochemical cells

their LHE is higher than that of transparent films as the optical path length is increased by the scattering.

Electron collection efficiency

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For complexes 1-3, the recombination of the injected electrons with the oxidized dyes, if anything, should be more efficient for 1 due to the higher surface concentration of

oxidized donors produced by the larger photocurrent density. Marcus has shown that, under a high-temperature approximation, the non-adiabatic electron-transfer rate between two electronic levels can be expressed by eqn. (4):

$$k_{\rm et} = \frac{2\pi |H_{\rm DA}|^2}{h(4\pi\lambda RT)^{1/2}} \exp \frac{-(\Delta G^0 + \lambda)^2}{4\lambda RT}$$
(4)

where $H_{\rm DA}$ is the electronic coupling matrix element, ΔG^0 is the free energy change for electron transfer process, and λ is the free energy of reorganization. We can estimate ΔG^0 for electron transfer from the conduction band edge to the oxidized dyes as $\Delta G_1^0 = -1.26 \text{ eV}$ for complex 1, $\Delta G_2^0 = -1.31 \text{ eV}$ for complex 2, $\Delta G_3^{0} = -1.40$ eV for complex 3 (here $E_{\rm CB} = -0.4$ V is assumed²¹). Reorganizational energies for the three complexes in this study are expected to be similar, and a typical value in related polypyridyl ruthenium complexes is 0.5 eV.36 These electron transfer processes should fall in the inverted region when $-\Delta G^0 > \lambda$. The larger the value of ΔG^0 , the lower the electron transfer rate. So, the rates of the electron transfer from the conduction band to the oxidized dyes, k_3 , for complexes 1-3 should be $k_{31} > k_{32} > k_{33}$ where k_{31} , k_{32} and k_{33} represent the rate of electron transfer of complexes 1, 2 and 3 respectively. Therefore, the molecular level similarity of the dyes and the relative values of k_3 allow us to conclude that η is not the cause of the differences in IPCE.

Electron injection quantum yield

Electron transfer from the excited dye to the TiO₂ surface will be an activationless process, if there is maximum overlap between the excited donor levels and the wide conduction band acceptor. The rate of activationless electron transfer from dye molecules to semiconductor surfaces is ultrafast and a lower limit of $k_2 > 10^9 \text{ s}^{-1}$ has been recently measured.³⁷ The quantum yield for electron injection (ϕ) is given by eqn. (5):

$$\phi = k_2 / (k_2 + k_r + k_{nr}) \tag{5}$$

where k_r and k_{nr} are the radiative and non-radiative rate constants for the excited dye and k_2 is the electron injection rate (Fig. 5). A lower ϕ would occur if radiative or nonradiative decay competes with electron injection. If the electron injection process is activated, the electron transfer rate will also decrease. As an approximation, electrochemical and spectroscopic data obtained in fluid solution can be used. The flat potential of TiO₂ film electrodes is -0.55 V (SCE) obtained from a Mott–Schottky curve which was determined by using a three-electrode system in MeCN. We employ the Rehm–

complex	ğhanaa ahaa ahaa ahaa ahaa ahaa ahaa aha	כ	ಯ, 1, 1, 1, 1, 1, 1, 1, 1, 1, 1, 1, 1, 1,	ሊ [†]	ಏсы	$\Delta G^0/{ m eV}$
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Weller equation [eqn. (6)]

$$\Delta G^{0} = E_{\rm ox} \,({\rm D}) - E_{\rm red} \,({\rm A}) - E^{0-0} - C \tag{6}$$

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As is well known for transition metal complexes, only the lowest excited state has a chance to emit luminescence and /or to live long enough to participate in bimolecular process. In the great majority of ruthenium(II) polypyridine complexes, the lowest excited state is a ³MLCT level (i.e., metal-to-ligand change transfer). Extensive investigation of ruthenium polypyridine complexes have given evidence for a single ligand localized excited state on the timescale of molecular vibrations with the excited electron localized on the ligand which is easier to reduce.³³ As mentioned above, for complexes 1, 2 and 3, the excited electron is mainly localized on ttp-PO₃H, phen and biquin, respectively in solution. When attached to the TiO₂ electrode, the excited electrons are still localized on biguin in complex 3 and on ttp-PO₃H in complex 1 because of the large π^* orbital energy differences between ttp-PO₃H cf. biquin and dmpy. The excited electron can be partly localized on ttp-PO₃H or phen in complex **2** because of the small π^* orbital energy differences between ttp-PO₃H and phen. Within the timescale of molecular vibrations, if the excited electron is not transferred in time to the TiO₂ conduction band, it will be deactivated undergoing radiative or non-radiative processes. Because the excited electron of complex 1 is localized on the ligand ttp-PO₃H attached to the TiO_2 electrode, it can be injected efficiently into the TiO₂ conduction bands leading to the best IPCE. For complex 2, lower IPCE is shown because only some of the excited electrons are localized on ttp-PO₃H. The excited electron of complex 3 is localized on biquin which is remote from the TiO₂ electrode, and its excited state deactivates mainly by the way of non-radiation and hence **3** exhibits the lowest sensitization to the TiO_2 electrode.

Conclusion

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